

# Topology and many-body physics in synthetic lattices



Alessio Celi

# Synthetic Hofstadter strips as minimal quantum Hall *experimental* systems



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- Synthetic lattice (**Extradimension**)
- Synthetic strip as minimal integer quantum Hall systems
  - Edge states in narrow strips
  - Topological response in narrow strips

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  - Meissner/Vortex phase (in analogy to type II superconductors)
    - Effect of the **dimerization**:
      - Reverse of chiral current (single particle)
      - Commensurate-Incommensurate transition (strong interactions)

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- Interacting flux ladder from ion chains *and a puzzle*
- Prospects and *philosophical view*

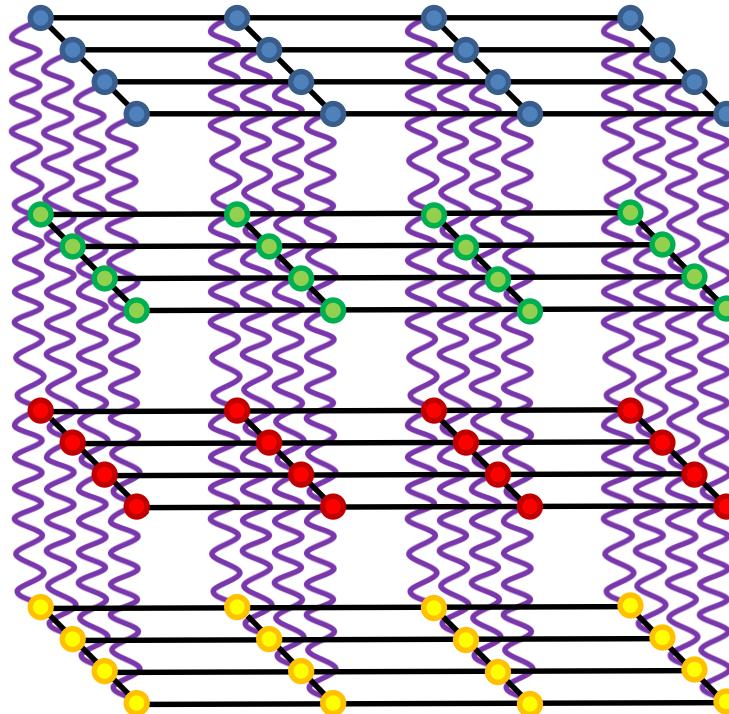
# Simulating an extra dimension

[Boada,AC,Latorre,Lewenstein, PRL 108, 133001 (2012)]

In optical lattices 1D-3D Hubbard model by tuning optical potential

And > 3D? In a lattice **Dimensionality**  $\equiv$  **Connectivity**

Atoms in D+1-lattice = coherently coupled atomic states in  $D$ -lattices

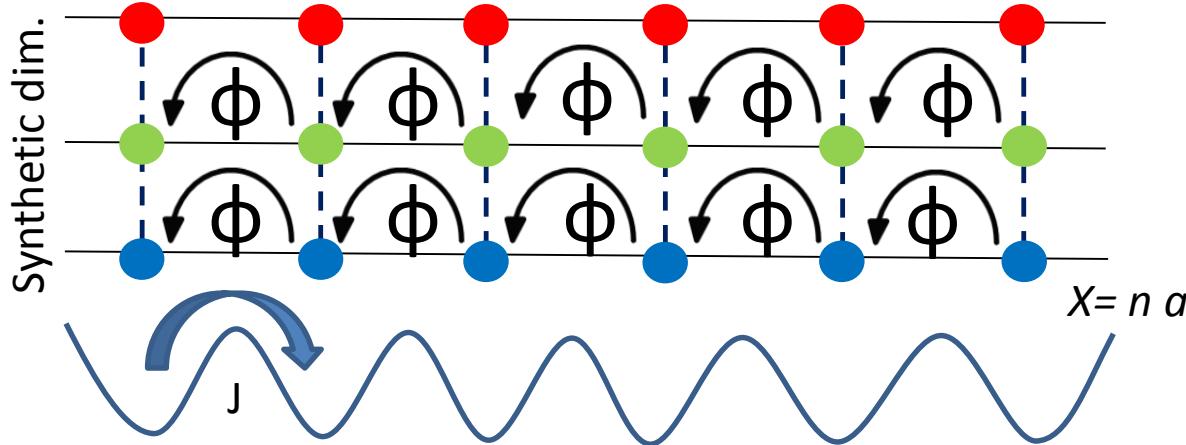


Not only spin states  
Momentum states  
Trap modes...

Not only atoms  
Cold molecules,  
Photonic crystal,  
Ring resonators...

# Synthetic gauge fields in synthetic dimension

[AC et al PRL 112, 043001 (2014)]



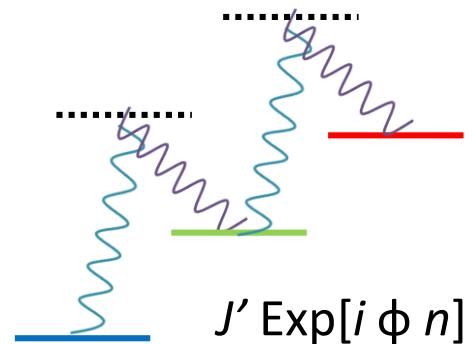
Constant magnetic flux  $\phi$ !

1d-lattice loaded e.g. with

$^{87}Rb$  ( $F=1, m=-1,0,1$ )

+

Raman dressing

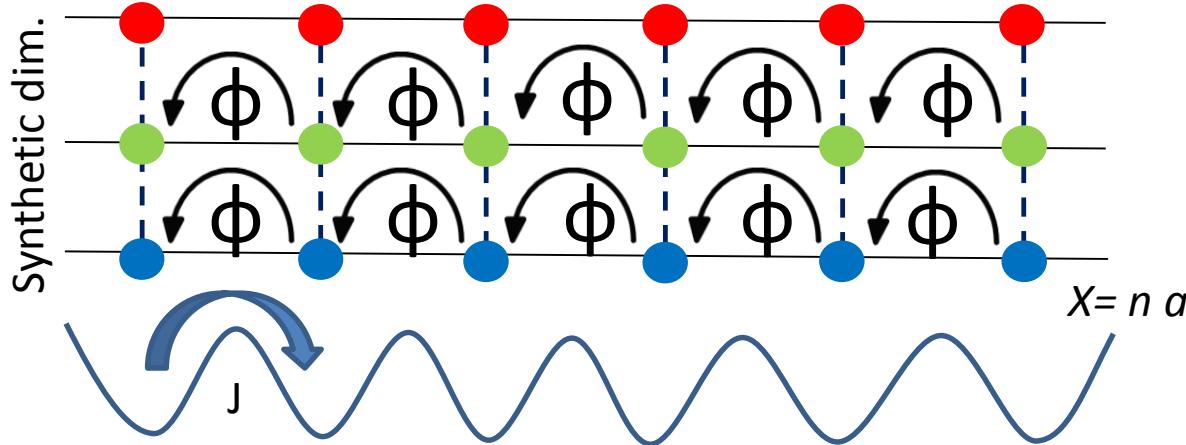


$$J' \exp[i\phi n]$$

Minimal instance of a quantum Hall system!

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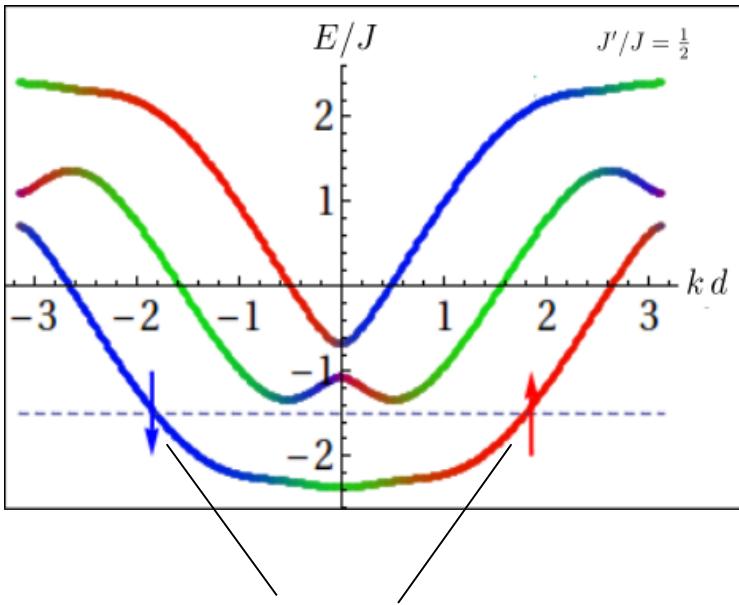
A schematic diagram showing a wavy blue line representing a synthetic dimension being dressed by a wavy purple line. A red horizontal line is also shown. Below the wavy lines is the expression  $J' \text{Exp}[i \phi n]$ .

Sharp Boundaries  $\rightarrow$  Edge currents (hard to get in real 2d lattice)  
signal of Topological nature of quantum Hall  
(bulk-boundary correspondence)

# Synthetic gauge fields in synthetic dimension

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## Spectrum



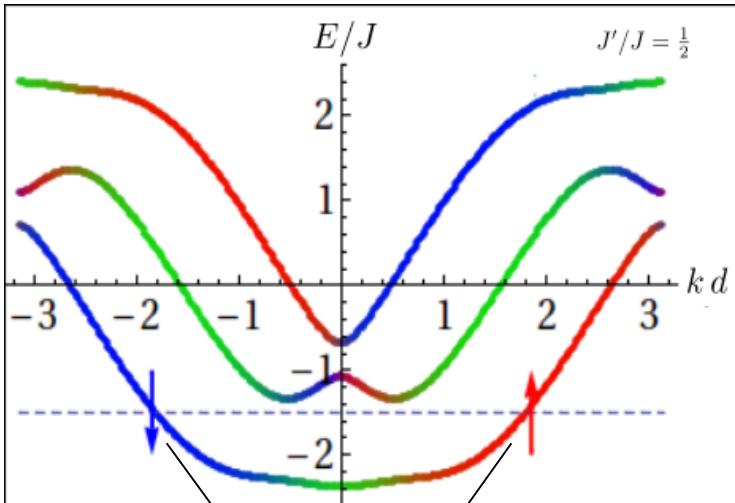
"Genuine" Edge states for small  $J'/J$ :

- live in the gap,
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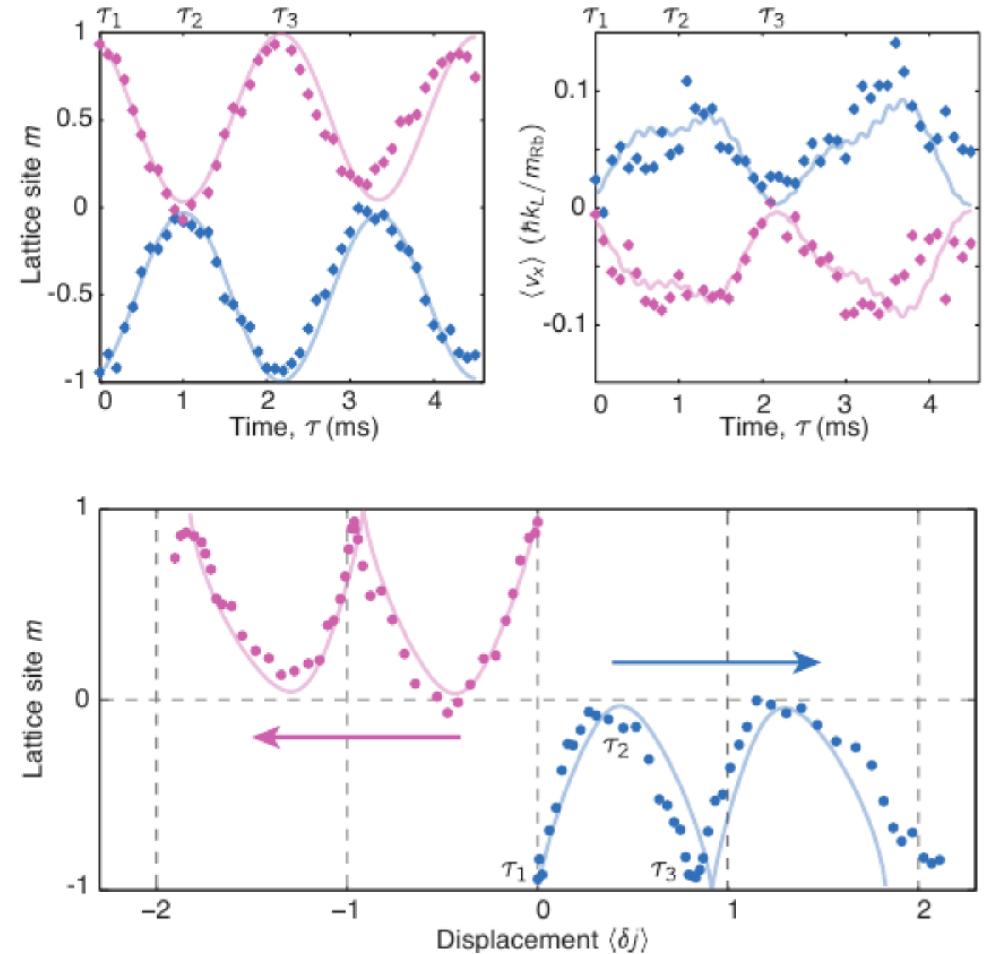
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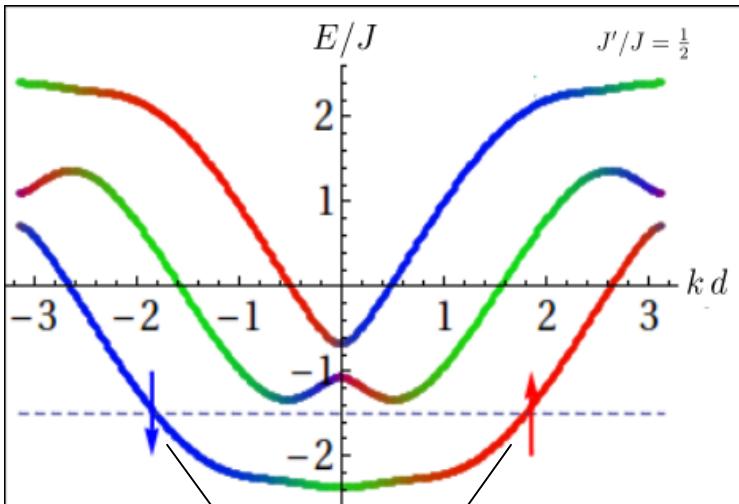
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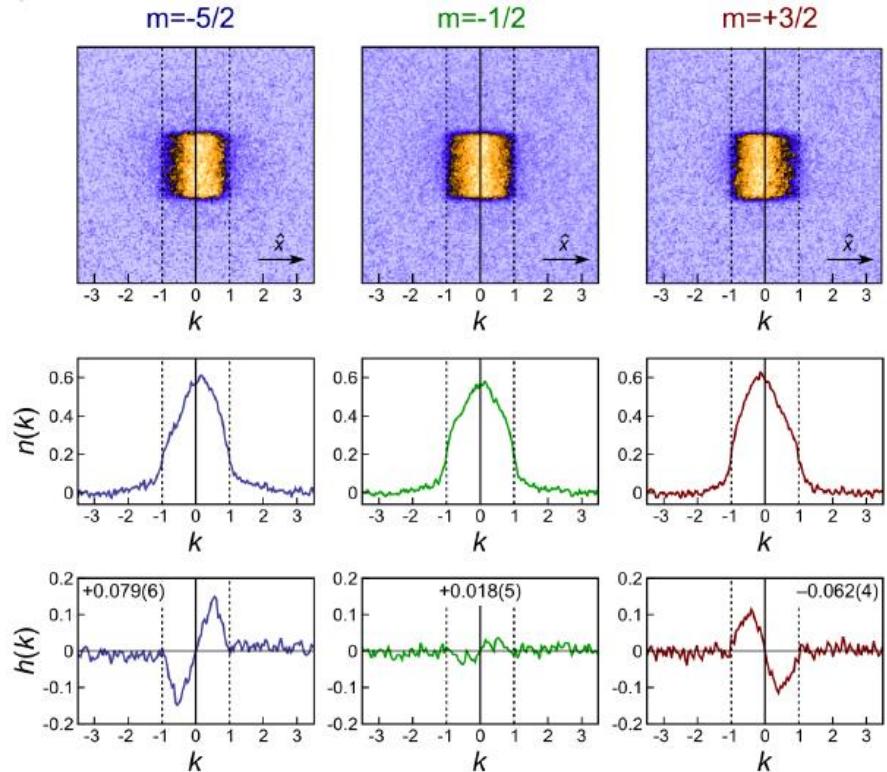
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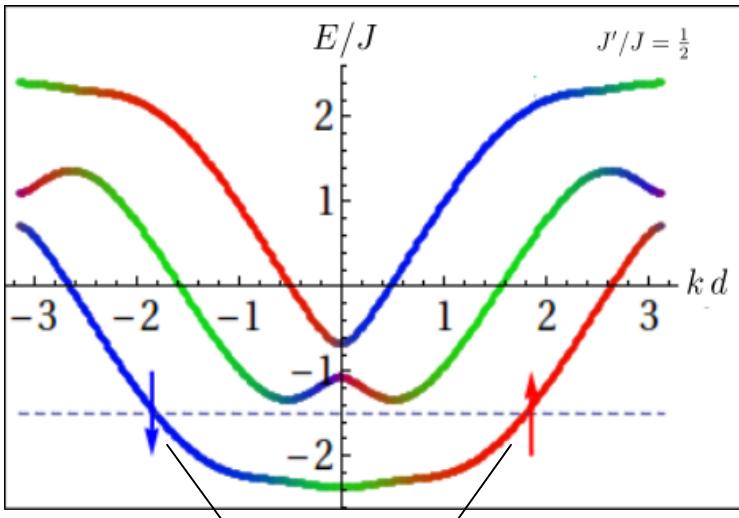
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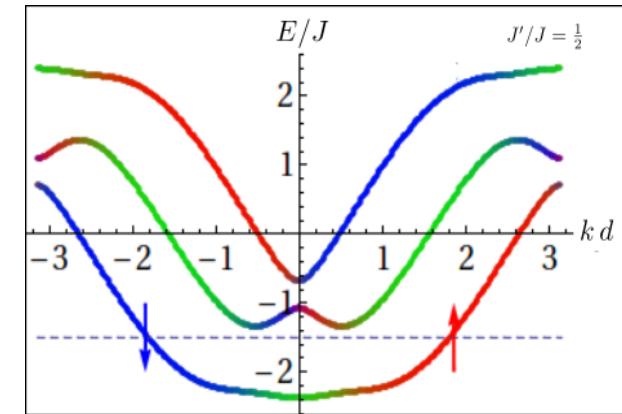
Also with clock states (ladder)

LENS: Livi et al. PRL 117, 220401 (2016)

JILA: Kolkowitz et al. Nature 542 66 (2017)

# Topology in narrow strips

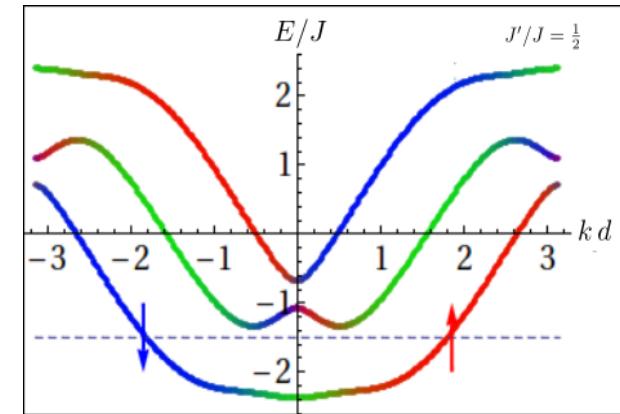
Narrow Hofstadter strips have edge states



What about the “bulk”?

# Topology in narrow strips

Narrow Hofstadter strips have edge states



What about the “bulk”?

How big should it be to display topological properties?

Is there some reminiscence of open/closed boundary correspondence?

Is Chern number defined? / Can we measure it?

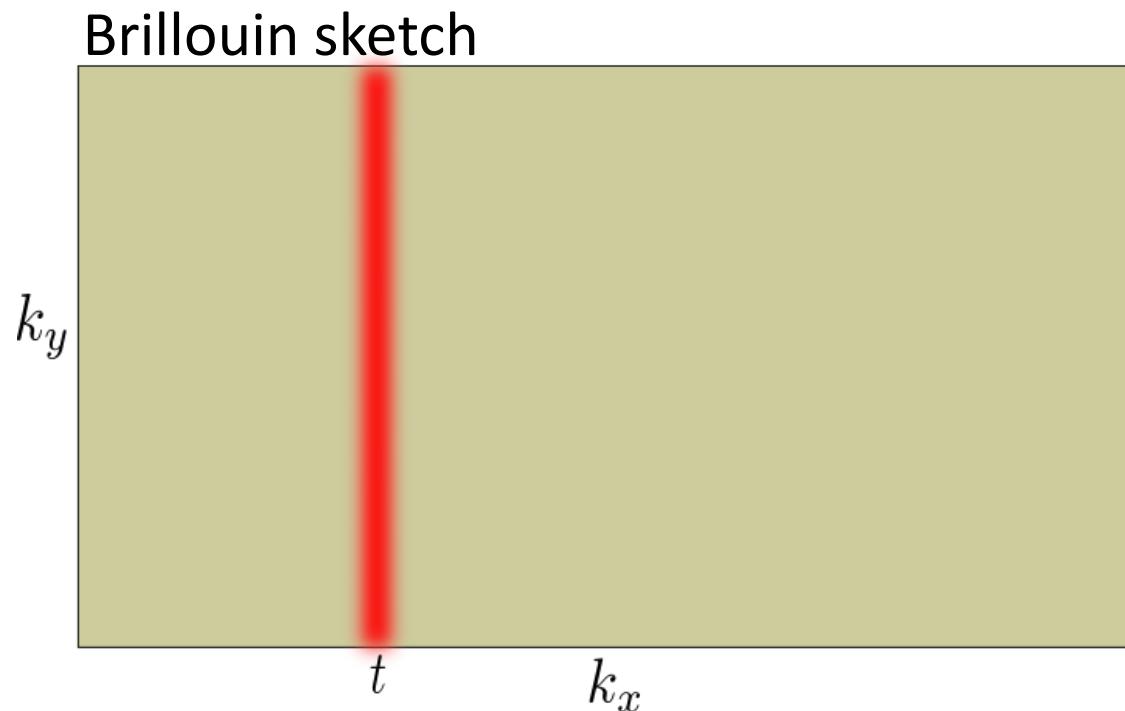
# Measuring Chern numbers in (narrow) Hofstadter strips, S.Mugel,...AC, SciPost 3, 012, (2017)

Pragmatic approach: measure transverse displacement to a force after  
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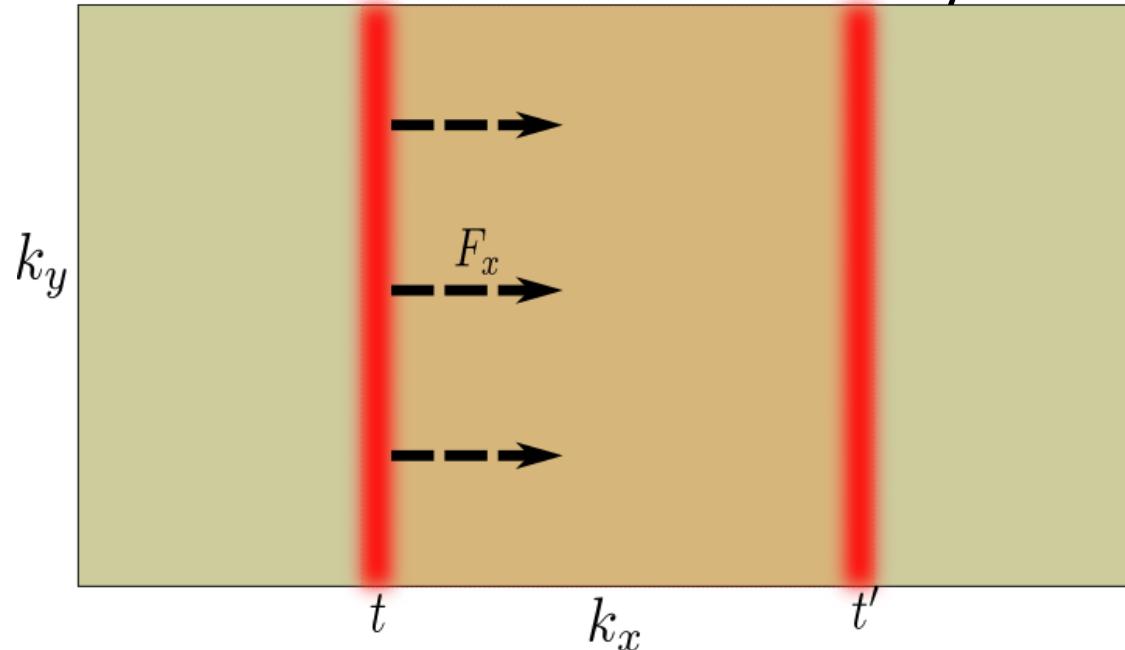
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Brillouin sketch of semiclassical dynamics



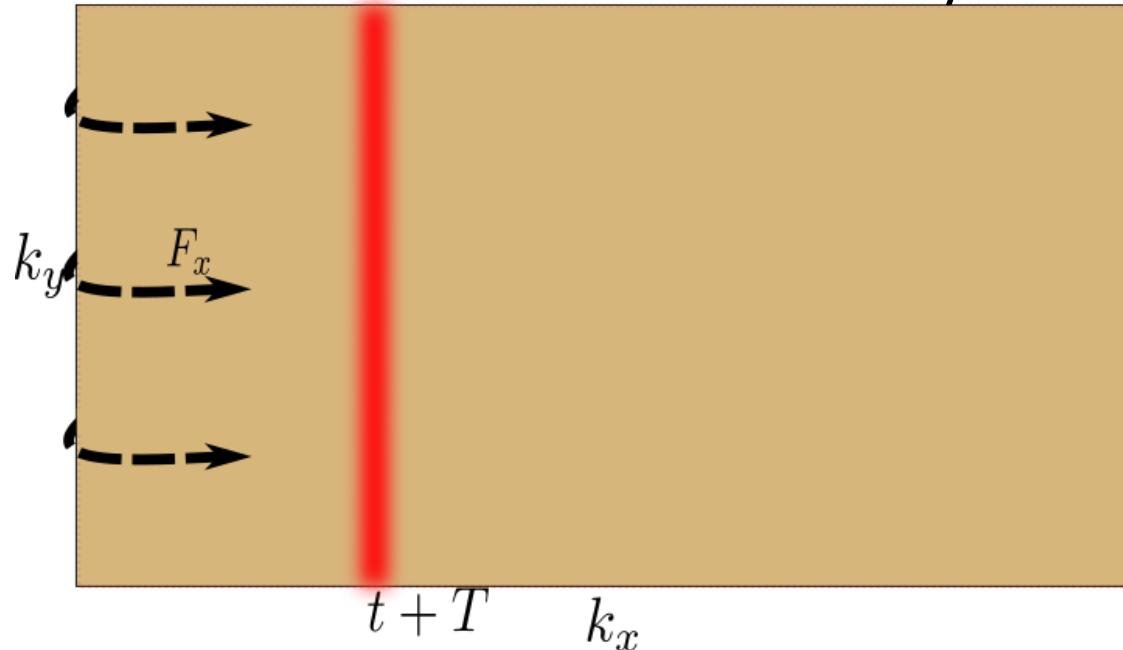
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- Apply a force along  $x$
- After a Bloch oscillation observe the displacement

Brillouin sketch of semiclassical dynamics



Displacement in  $y$   
due to anomalous  
velocity!

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Pragmatic approach: measure transverse displacement to a force after  
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*In formulae:* semiclassical approach

$$\mathbf{k}(t) = \mathbf{k}_0 + \frac{t}{\hbar d} F_x \mathbf{e}_x \quad \longrightarrow \quad \mathbf{v}(\mathbf{k}) = \frac{1}{\hbar} \partial_{\mathbf{k}} E(\mathbf{k}) + \frac{F_x}{\hbar d} \mathcal{F}(\mathbf{k}) \mathbf{e}_y$$

$$|\psi(\mathbf{k})|^2 \sim \frac{1}{A_{BZ}} \delta(k_x - k_x(t))$$

Wave packet:

$$\langle \mathbf{r}(T) - \mathbf{r}(0) \rangle = \int_0^T \langle \mathbf{v}(t) \rangle dt = \frac{\hbar d}{|F_x|} \int_{BZ} \mathbf{v}(\mathbf{k}) = \text{sgn}(F_x) \mathcal{C} d \mathbf{e}_y$$

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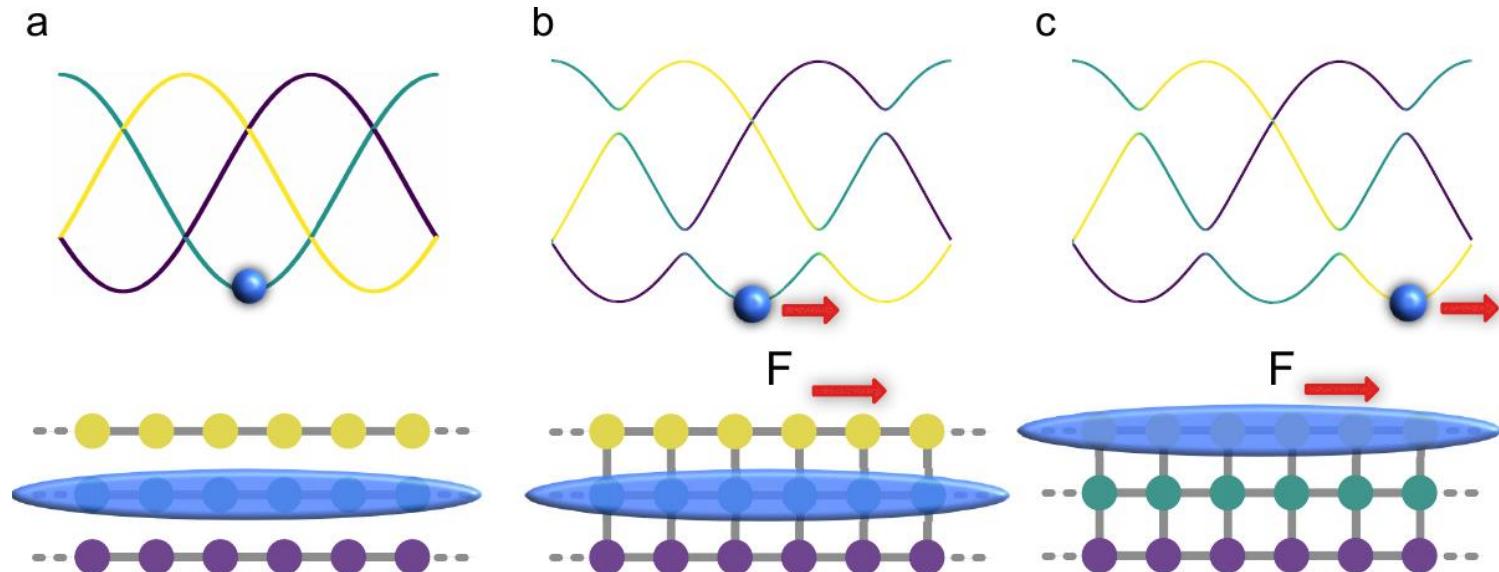
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Applicable also to strips until we don't reach the boundary...

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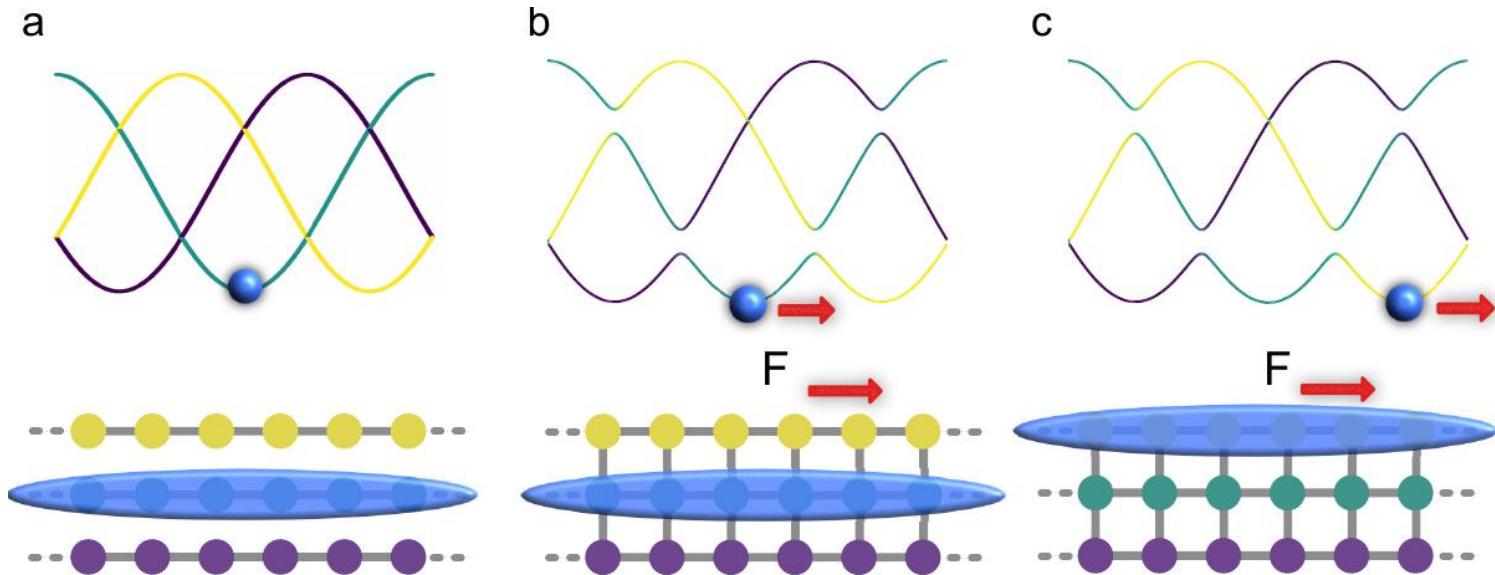
Scheme:



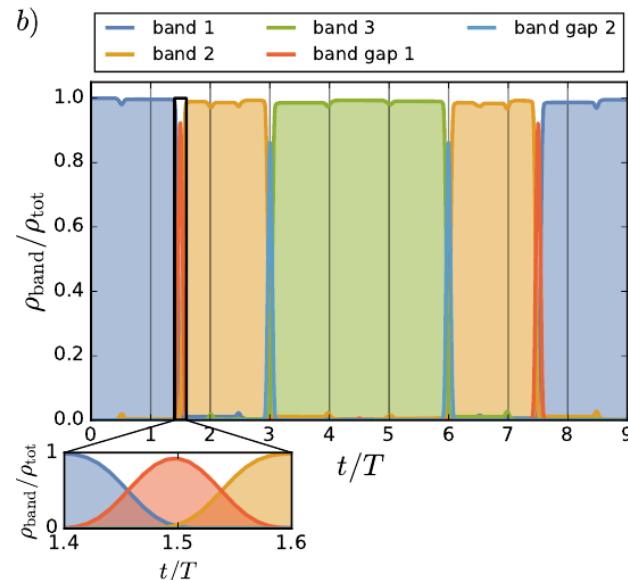
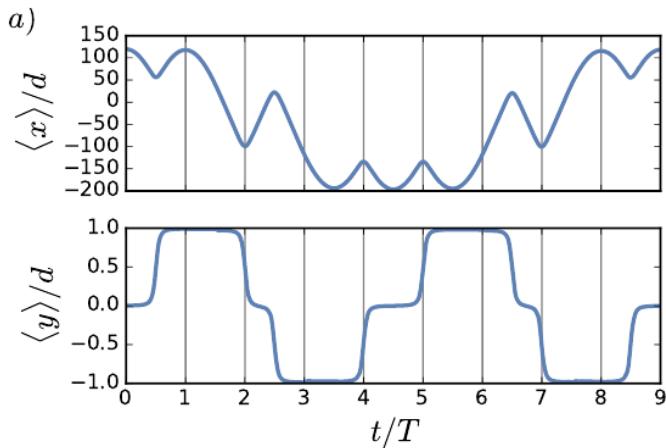
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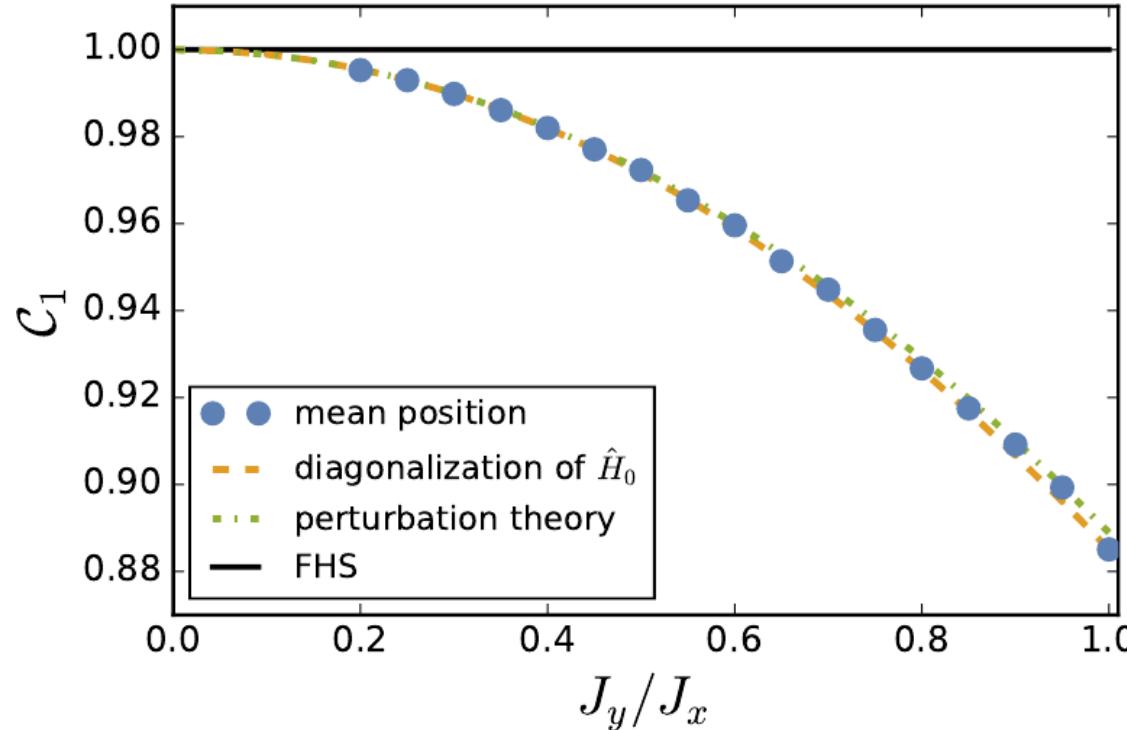
Results:  $J_y = \frac{1}{5}J_x$ ,  $\Phi = \frac{2\pi}{3}$   $N_y = 3$



# Measuring Chern numbers in (narrow) Hofstadter strips, S.Mugel,...AC, SciPost 3, 012, (2017)

Why does it work? **Perturbative argument** as for the edge states:

- Gap linear in  $J_y/J_x$
- Hybridization spin states (spreading in  $y$ ) quadratic in  $J_y/J_x$

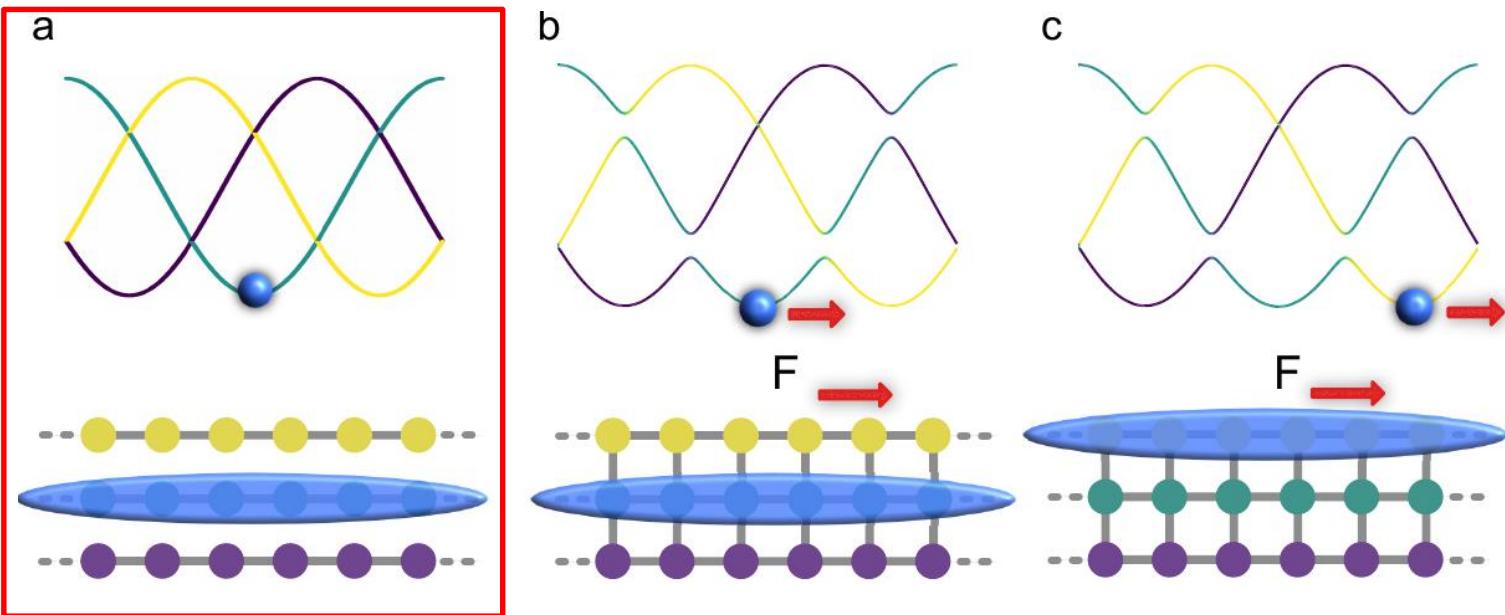


Quadratic degradation of the measurement

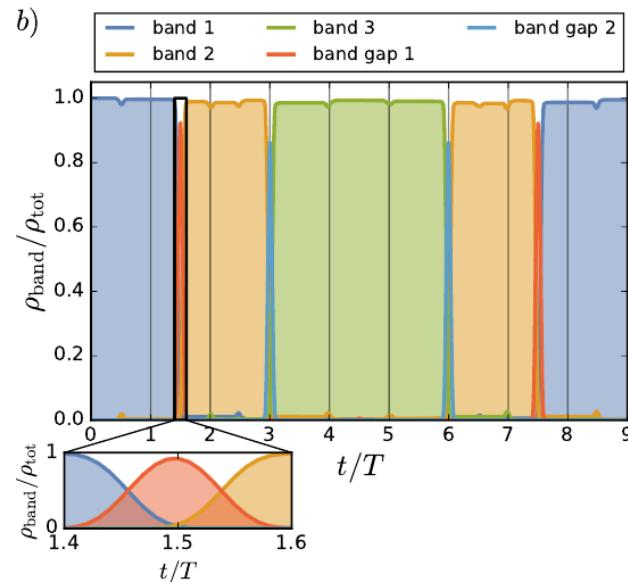
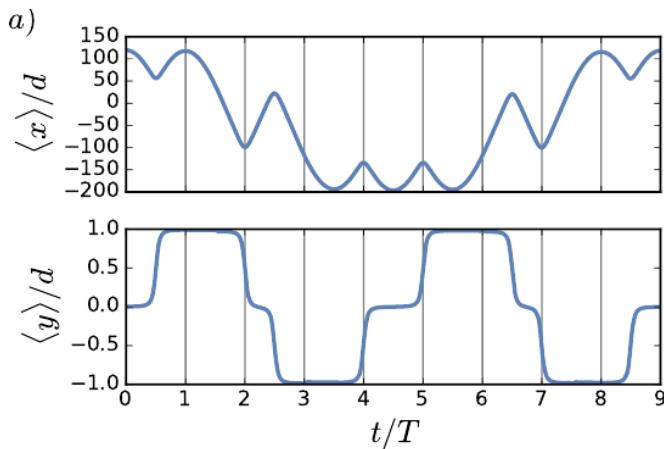
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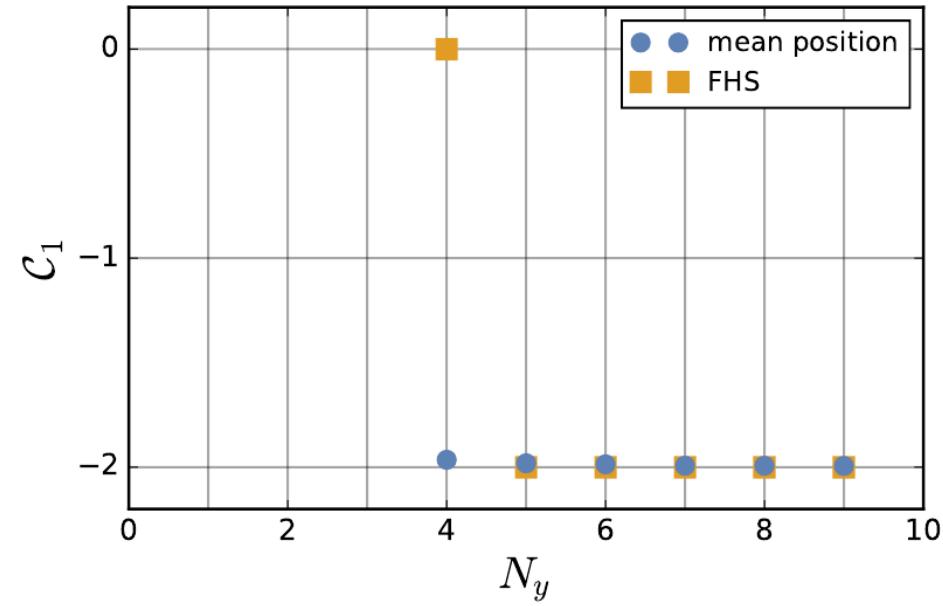


# Measuring Chern numbers in (narrow) Hofstadter strips, S.Mugel,...AC, SciPost 3, 012, (2017)

Higher  $\mathcal{C}$  possible for  $N_y \geq \mathcal{C} + 2$

Ex:  $\Phi = \frac{4\pi}{5} \rightarrow \mathcal{C}_1 = -2$

“Better” than Fukui-Hatsugai-Suzuki algorithm J. Phys. Soc. Jpn. (2005)



Robust to disorder (< gap) and typical harmonic confinement

Interactions? Gap small (although may hold thought adiabatic argument, see later)

# Synthetic lattices in interaction

- Interesting route to interaction -> Fractional QH effect?!  
-> many-body localization?!
- No heating expected
- Peculiarity: Interactions are naturally long range  
in the synthetic dimension
- Quasi 1D approach to 2D interesting both  
theoretically & practically

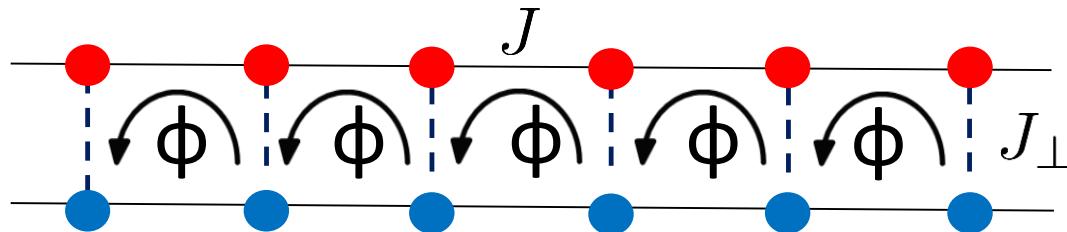
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Many studies: Meissner-vortex and commensurable incommensurable transitions,  
Fractional pumping, Laughlin like states, pseudo Majorana...

Here: effect of dimerization on synthetic Hofstadter ladder

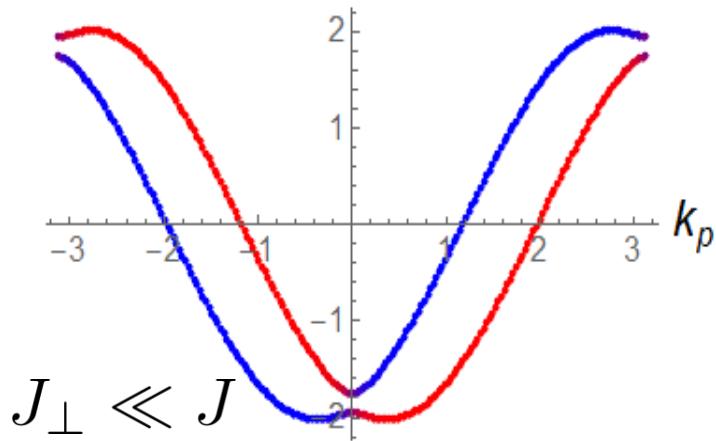
# Meissner/Vortex phase in flux ladder



No interactions: real = synthetic ladder

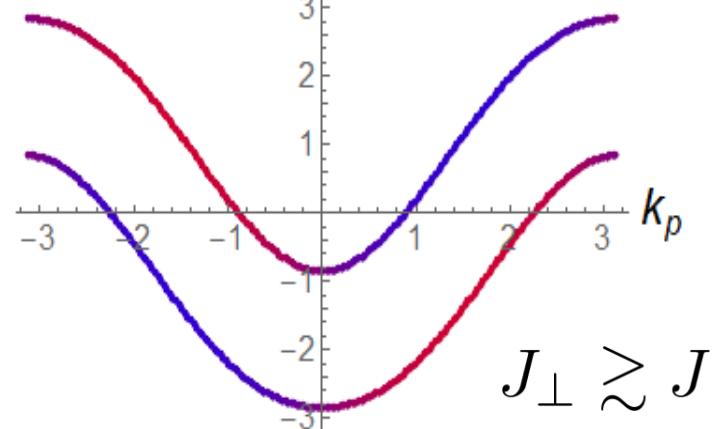
Weak interleg (Raman) coupling:

$$2 \text{ minima, } k_m \sim \pm \frac{\phi}{2}$$



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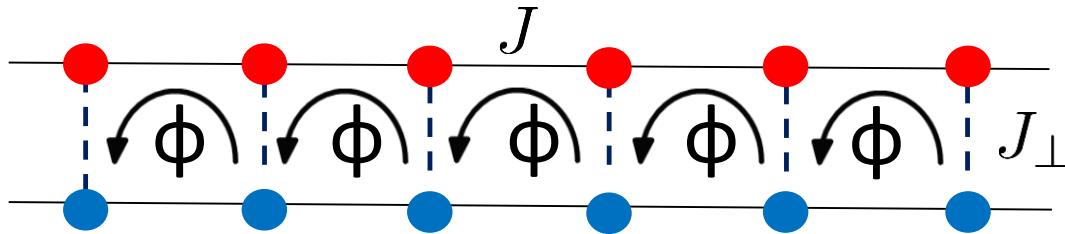
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[Orignac,Giamarchi, PRB 2001] Analogous to type II, also in presence interactions

Real ladder experiment [Atala et, Nature Phys. 2014]

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## Observables

$$J_c(j, m) = i \langle \hat{a}_{j+1, m}^\dagger \hat{a}_{j, m} \rangle + H.c.$$

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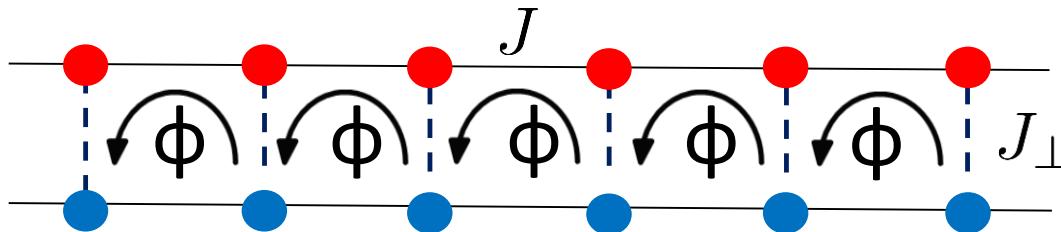
$$J_\perp \ll J$$

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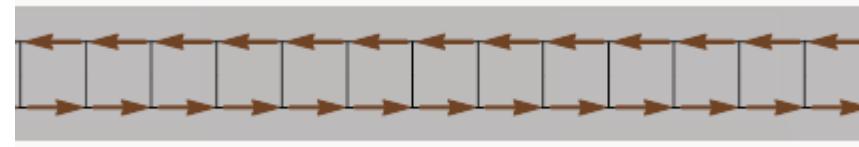
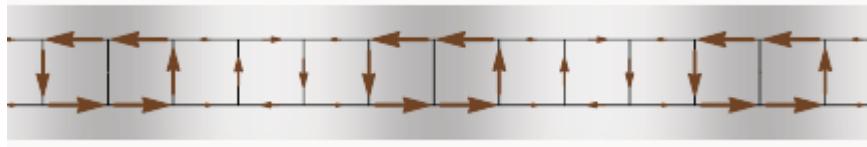
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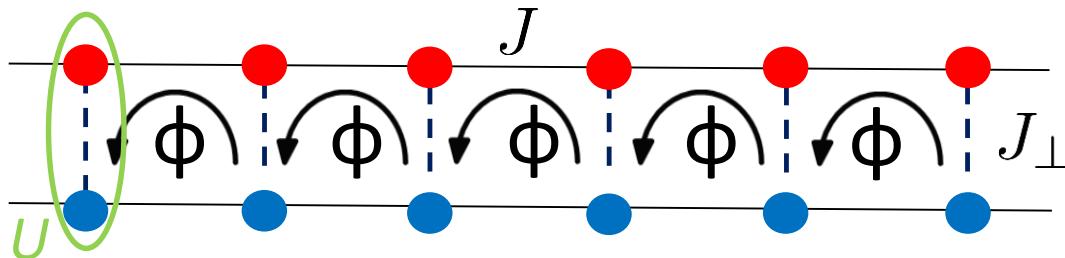
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Effect of interactions: suppress vortex phase

Real ladder: vortex phase survives in the hard-core limit for  $\phi$  large

more phases at  $U \neq \infty$

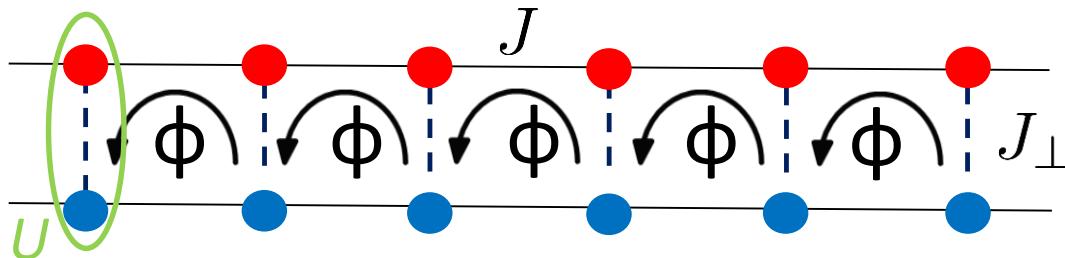
see [Petrescu, Le Hur, PRL 2013]  
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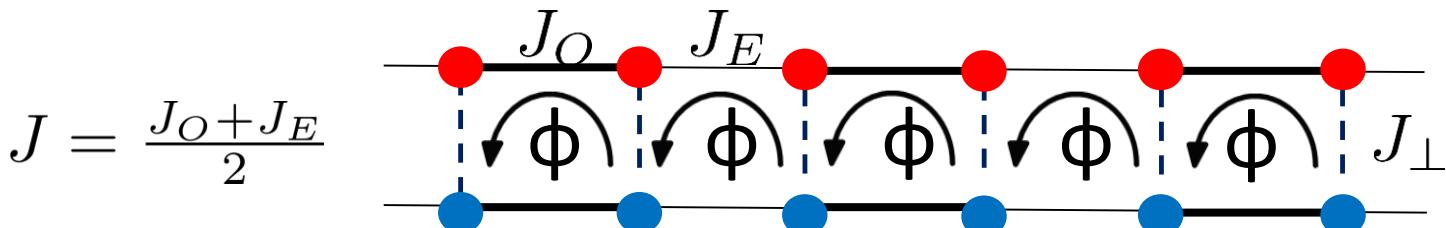
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Idea: enhance vortex phase by dimerizing the lattice (“easy” exp. handle)

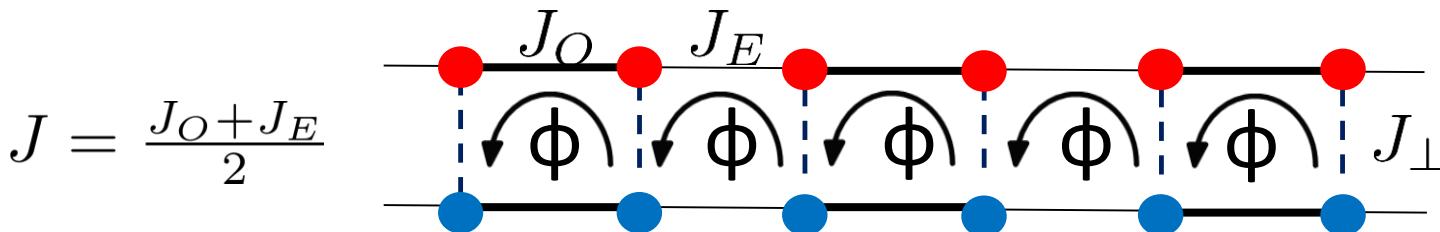
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$$J = \frac{J_O + J_E}{2}$$

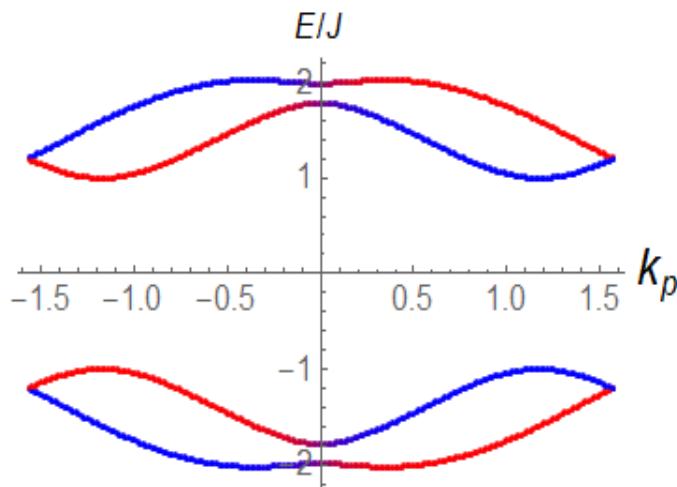
Effect of dimerization: new handle  $\Delta = \frac{J_O - J_E}{J_O + J_E}$

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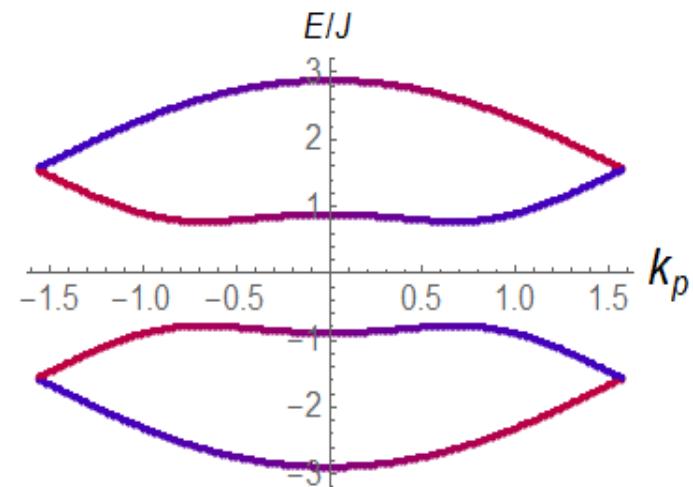
No interactions: 4 bands



$$J_{\perp} \ll J$$

$$\Delta = 0.5$$

Bands  
deform  
& mix

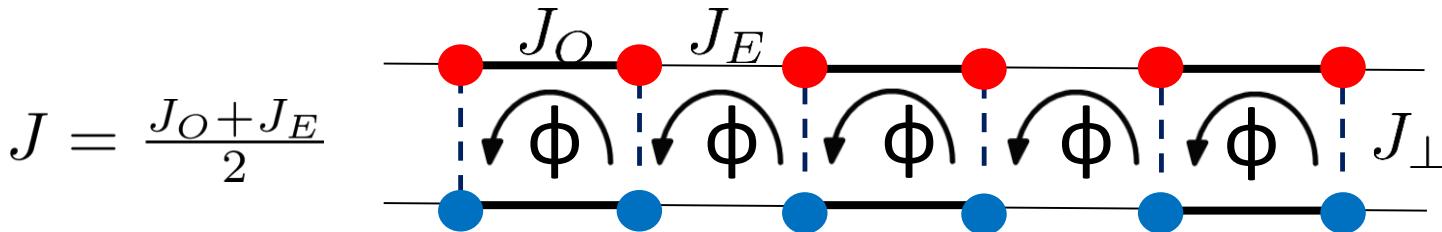


$$J_{\perp} \gtrsim J$$

Minima separate: dimerization enhances vortex phase!

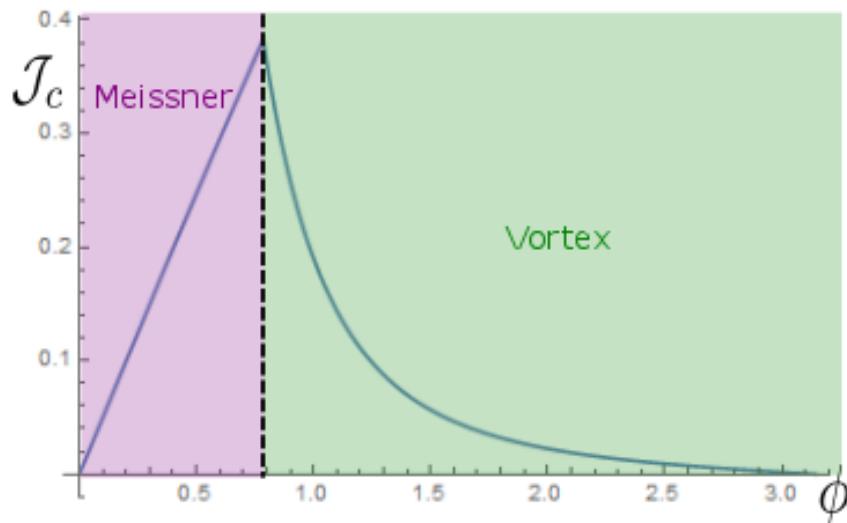
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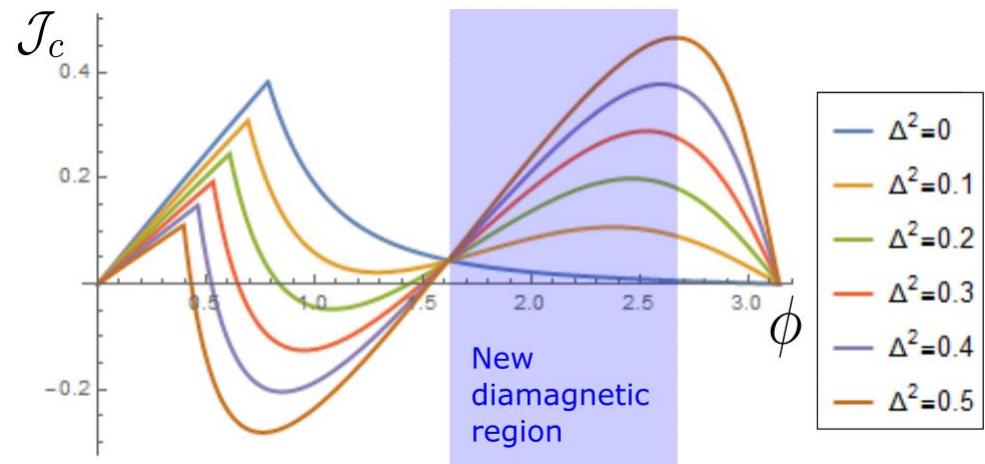


Effect of dimerization: new handle  $\Delta = \frac{J_O - J_E}{J_O + J_E}$

No interactions: Reverse of chiral current



$$\Delta = 0$$

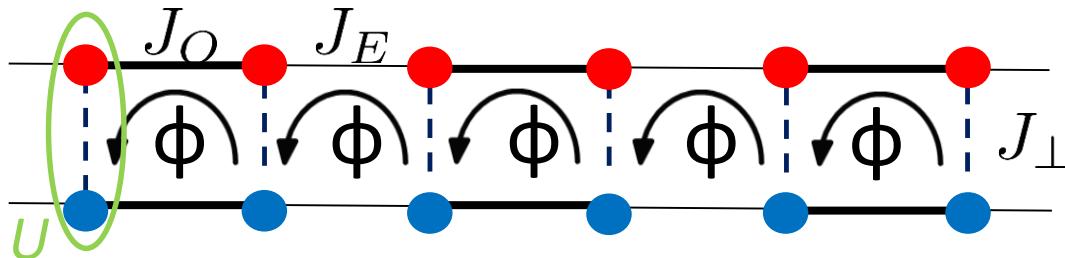


$$\Delta \neq 0$$

Current behavior confirms vortex enhancement!

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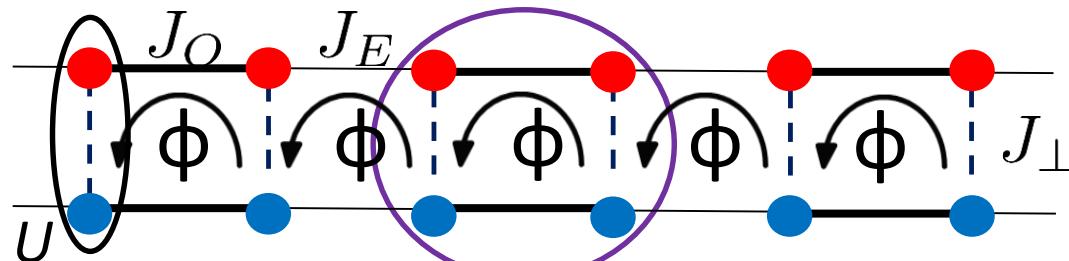


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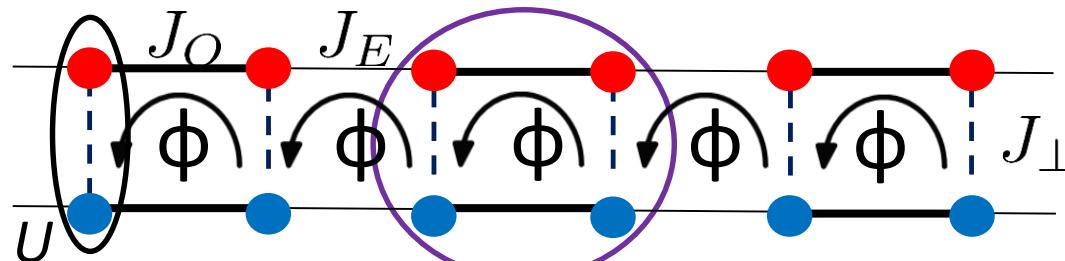
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$J_E \ll J_O$  9 states per plaquette

	1 $n=0,$	4 $n=1,$	4 $n=2$
Spectrum plaquette	0	$\pm J_{\perp} \sqrt{1 + \left(\frac{J_O}{J_{\perp}}\right)^2 - 2 \cos \frac{\phi}{2} \left(\frac{J_O}{J_{\perp}}\right)}$	$\pm 2J_{\perp}$
		$\pm J_{\perp} \sqrt{1 + \left(\frac{J_O}{J_{\perp}}\right)^2 + 2 \cos \frac{\phi}{2} \left(\frac{J_O}{J_{\perp}}\right)}$	$\pm 0$

# Vortex Nesting and Melting in synthetic ladders, E. Tirrito, R. Citro, M. Lewenstein, AC *in progress*

$$J = \frac{J_O + J_E}{2}$$



Effect of dimerization: new handle  $\Delta = \frac{J_O - J_E}{J_O + J_E}$

Interactions:  $U \rightarrow \infty$  3 states per rung

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	1 $n=0,$	4 $n=1,$	4 $n=2$
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$$J_{\perp} \gtrsim J_O$$

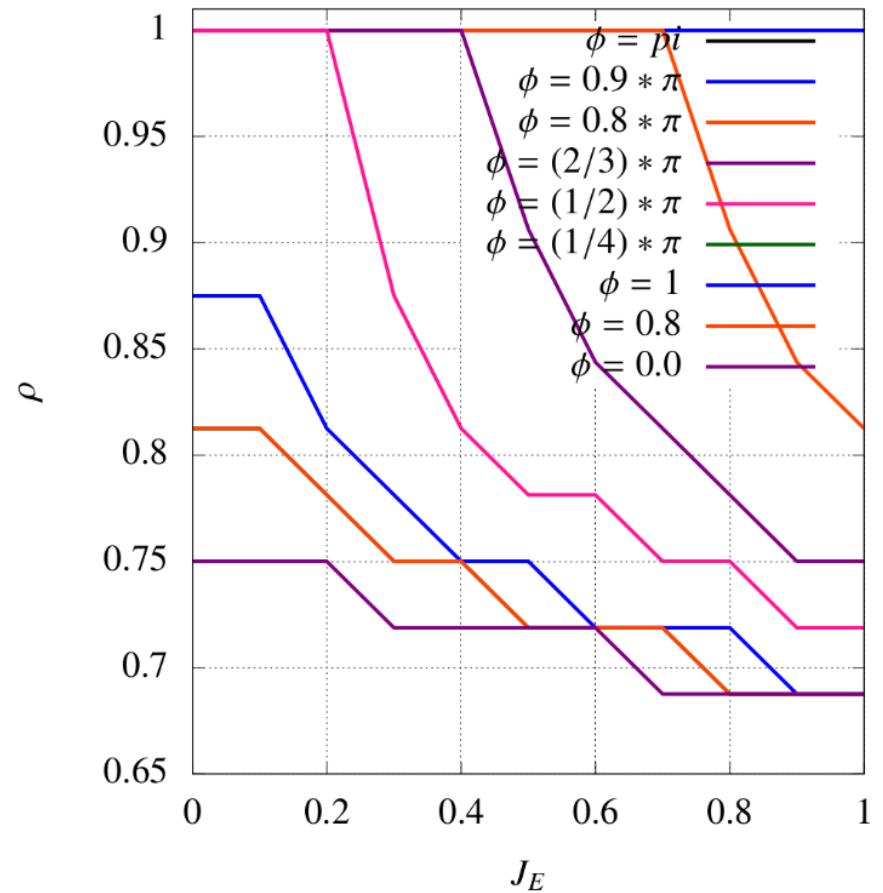
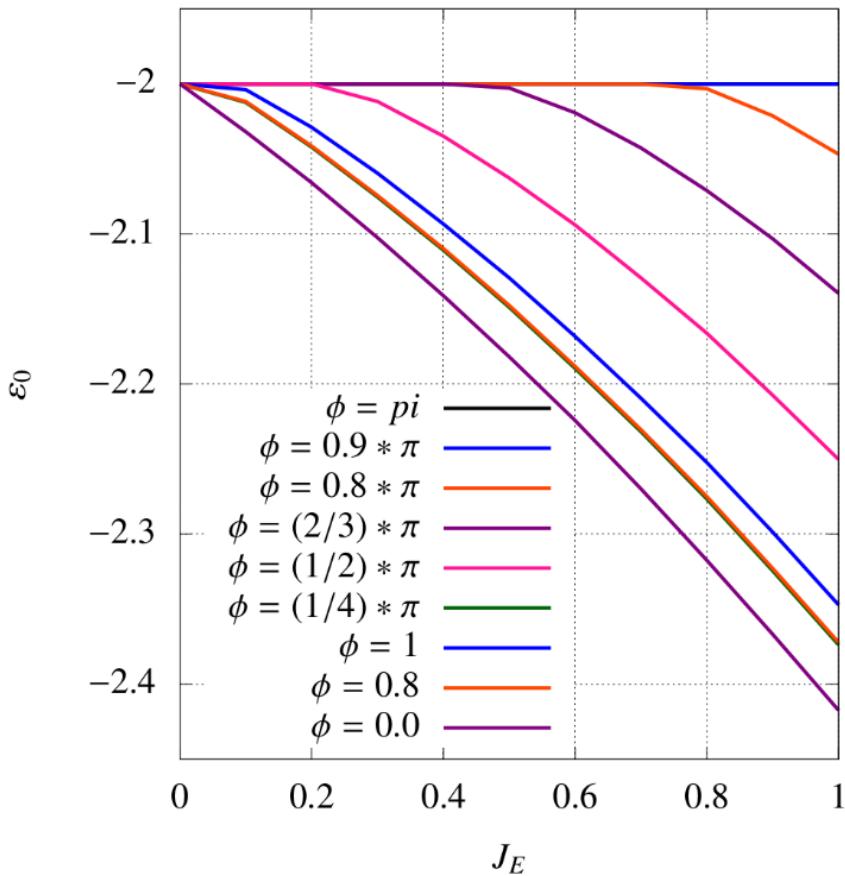
$$J_{\perp} < J_O$$

Plaquette in  $n=2$   $\longrightarrow$  Band insulator  
 Plaquette in  $n=1$   $\longrightarrow$  Imprinted vortex

# Vortex Nesting and Melting in synthetic ladders, E. Tirrito, R. Citro, M. Lewenstein, AC *in progress*

DMRG calculations confirm perturbative expectations

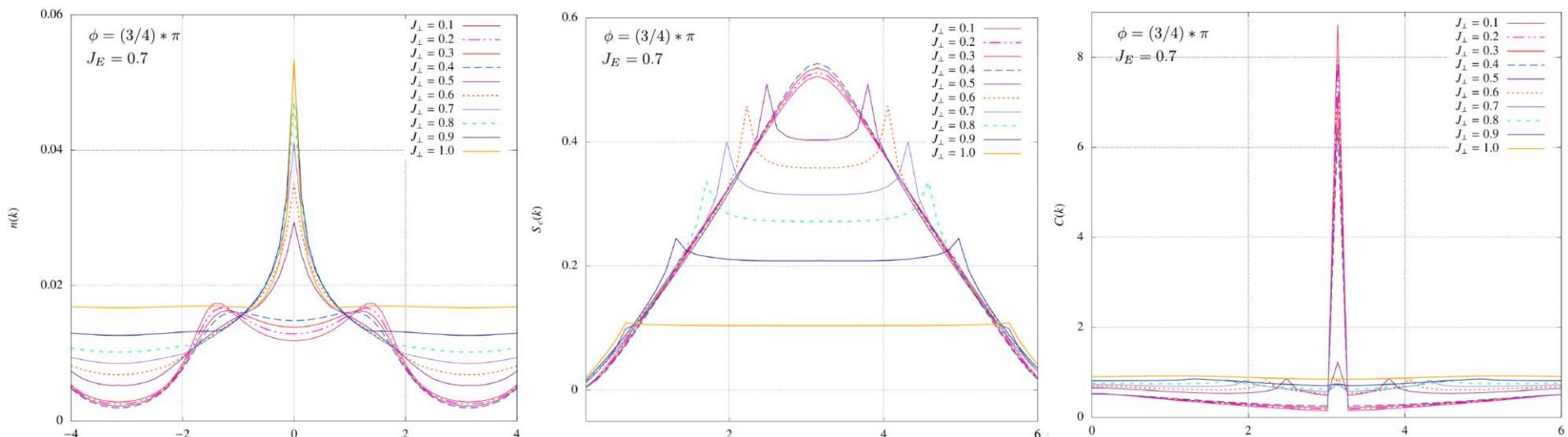
Ex.  $J_{\perp} = J_O = 1$



Phase diagram through calculation of currents and structure factors

# Vortex Nesting and Melting in synthetic ladders, E. Tirrito, R. Citro, M. Lewenstein, AC *in progress*

Ex.  $\phi = 3\pi/4$



0



$J_E$



$J_O$

Imprinted vortex → Melted vortex → Meissner charge density wave → Meissner

Similar to attractive interleg interaction, [Orignac et al, PRB 96, 014518 (2017)]

Evidence of commensurate-incommensurate transition

# *Further steps*

- No hard-core boson limit: bosons different fermions
- Study the accessible experimental parameters
- Search for “visible” Laughlin-like states in such regimes  
cf. [Calvanese et al, PRX 7, 021033 (2017)], [Petrescu et al, PRB 96, 014524 (2017)]
- ...*Toy model for many-body localization?*

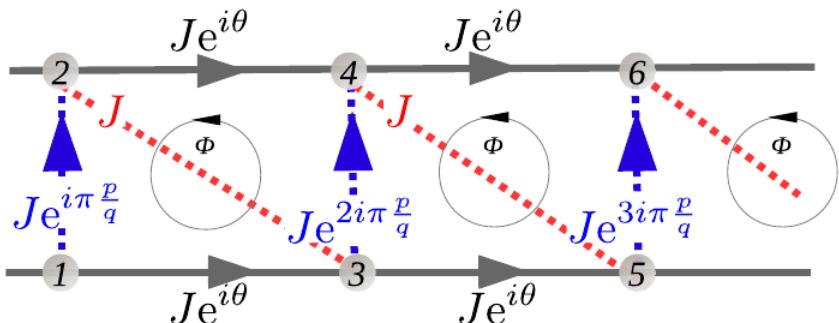
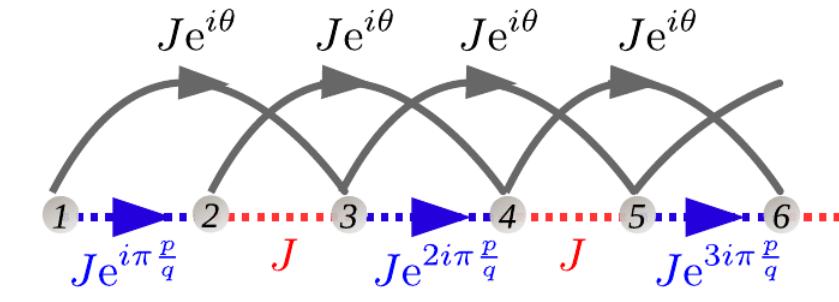
# Alternative route to synthetic interacting ladders... long range interactions!

PHYSICAL REVIEW A **91**, 063612 (2015)

## Synthetic magnetic fluxes and topological order in one-dimensional spin systems

Tobias Graß,<sup>1</sup> Christine Muschik,<sup>1,2,3</sup> Alessio Celi,<sup>1</sup> Ravindra W. Chhajlany,<sup>1,4</sup> and Maciej Lewenstein<sup>1,5</sup>

<sup>1</sup>ICFO-Institut de Ciències Fotòniques, Av. Carl Friedrich Gauss 3, 08860 Barcelona, Spain



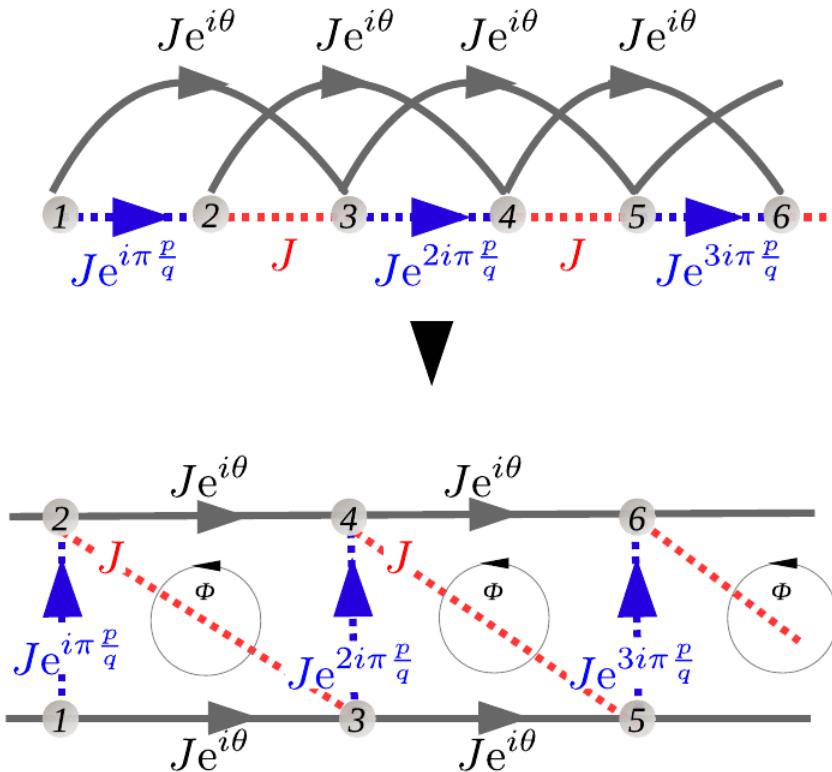
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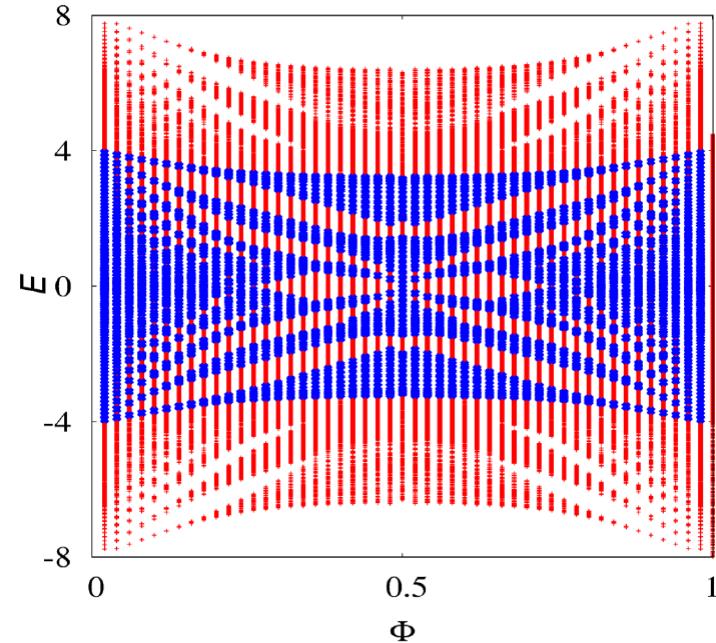
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How? E.g. ion analogue simulation

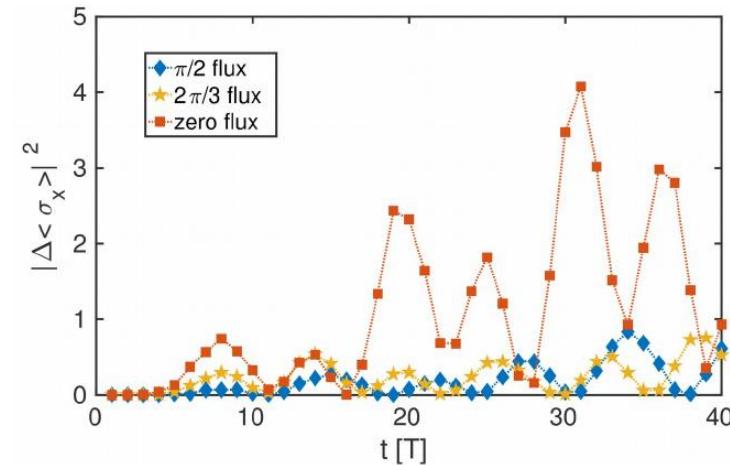
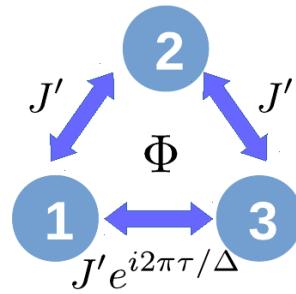
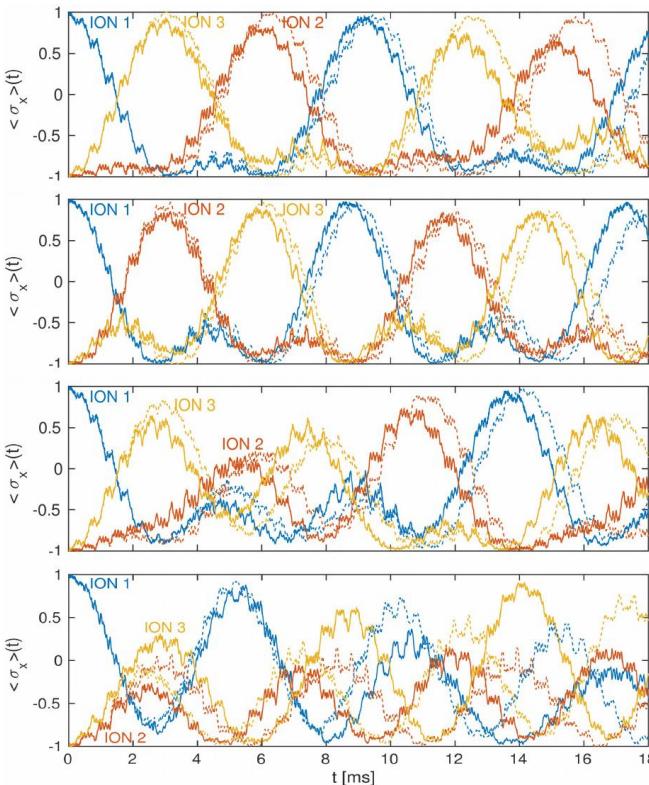


# Alternative route to synthetic interacting ladders... long range interactions!

Driven Dicke model with “fluxes” seems robust to photon heating!

$$H_0(t) = \sum_m \hbar\omega_m a_m^\dagger a_m + \sum_{i,m} \hbar\Omega_i \eta_{i,m} (a_m + a_m^\dagger) \sigma_i^x \sin(\omega t) + \sum_i B_i(t) \sigma_i^z$$

Special instance: Triangle! [T. Grass, AC, G. Pagano, M. Lewenstein, arXiv:1708.01882]



Deviation from the effective spin model are suppressed also for strong driving in presence of fluxes! *Topological protection?*

# *Summary*

- Synthetic edge state in synthetic Hofstadter strips
- “Bulk topology” in synthetic Hofstadter strips
- Effect of dimerization in synthetic Hofstadter ladder w/o interactions
- Synthetic fluxes in driven ion chain

# *A bit of philosophy...*

Synthetic lattices: practical tool for reshaping degrees of freedom in experimentally convenient way...

Same Hamiltonian = Same model

Useful approach for questioning conventions, ease detection, suggest new models...

# *A bit of philosophy...*

Synthetic lattices: practical tool for reshaping degrees of freedom in experimentally convenient way...

Same Hamiltonian = Same model

Useful approach for questioning conventions, ease detection, suggest new models...

Pushing it further...

Same time evolution = Same model

Synthetic lattices naturally combine with the Floquet approach...

Realization of quantum simulation paradigm

# *“Extradimensional” collaborators*



J.I. Latorre



O. Boada



M. Lewenstein

# *“Extradimensional” collaborators*



J.I. Latorre



M. Lewenstein



T. Grass

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